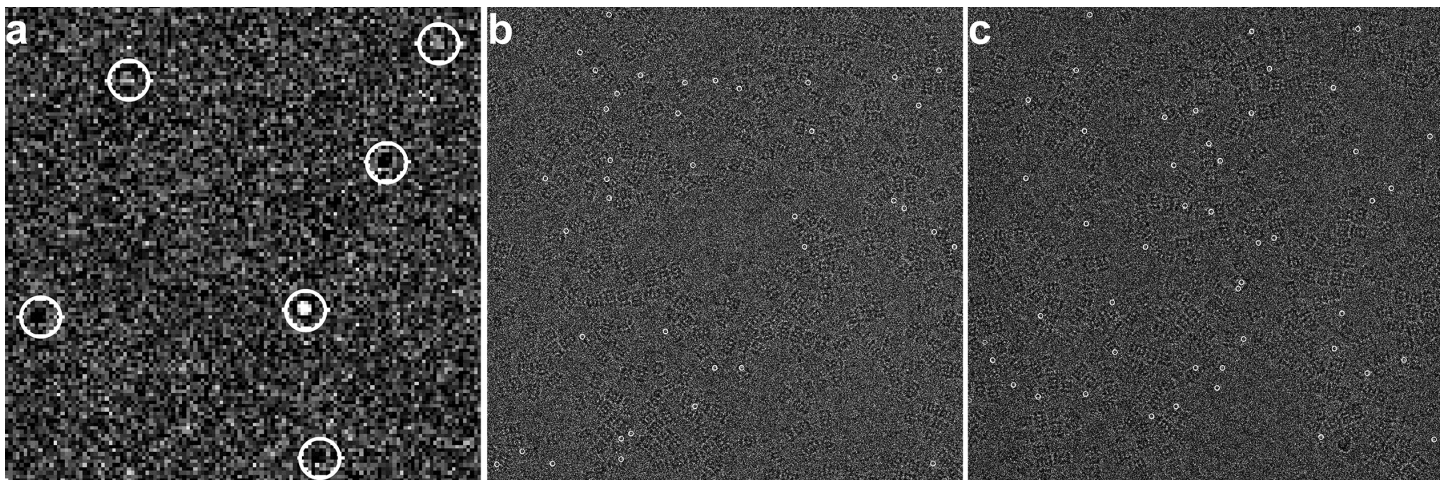


Supplementary Figure 1

A doming model describes the motion of frozen hydrated samples induced by the high-energy electron beam.

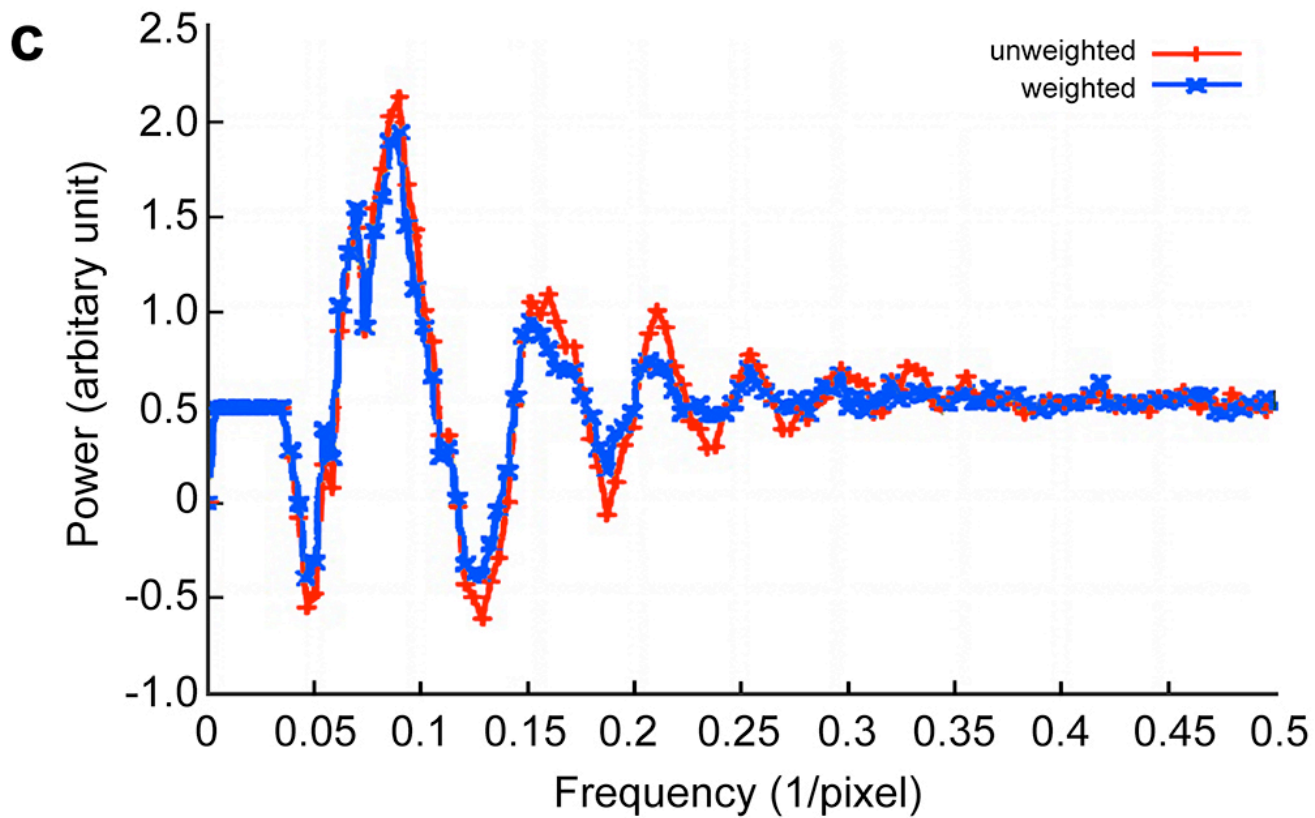
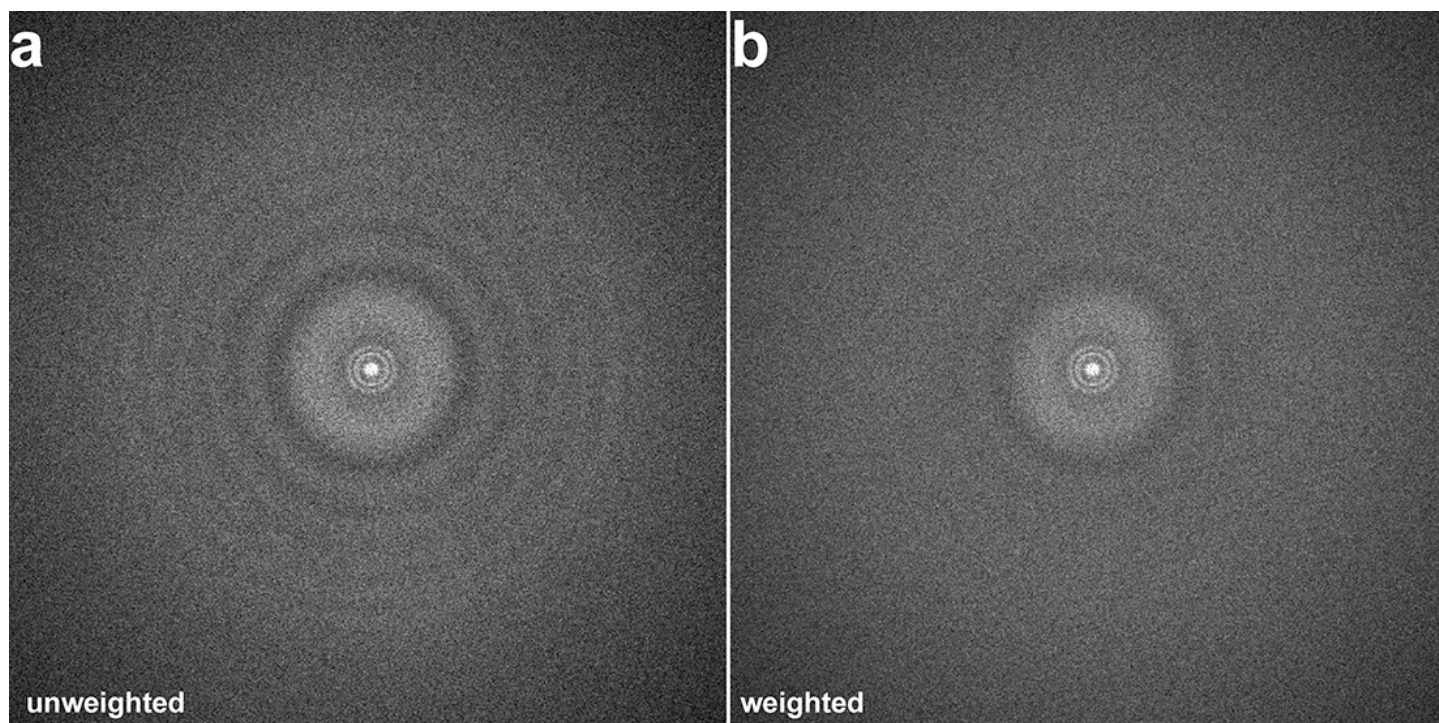
(a) Traces of the projected motion measured at three different tilt angles (α -angle) extracted from a dose fractionated tilt series acquired on a frozen hydrated specimen of a wild-type *Drosophila melanogaster* γ -Tubulin Ring Complex. The black arrow in the plot indicates the starting position of the motion. (b) Image of frozen hydrated archaeal 20S proteasome overlaid with the traces of global motion based upon whole frame alignment (long trace originated from the center of image) and each patch determined from MotionCor2. The whole frame is divided into 5×5 patches, and traces of each patch are determined individually. The traces of the global and local motions are exaggerated on the image by a factor of 145. For perspective, the accumulated global motion is ~ 11 Å.



Supplementary Figure 2

Defect pixel detection in MotionCor2

(a) An example showing defective pixels (marked with white circles) in an image captured with the K2 Summit camera. (b) and (c) Two different images collected one after the other show pixel defects (marked with white circles) in different locations.

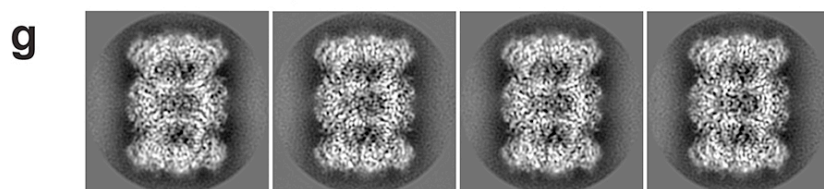
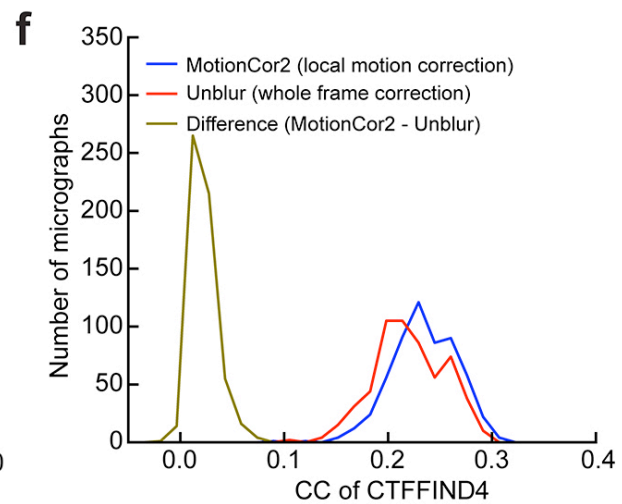
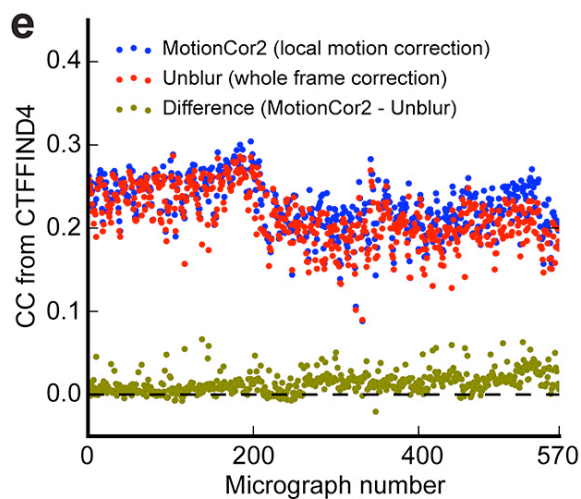
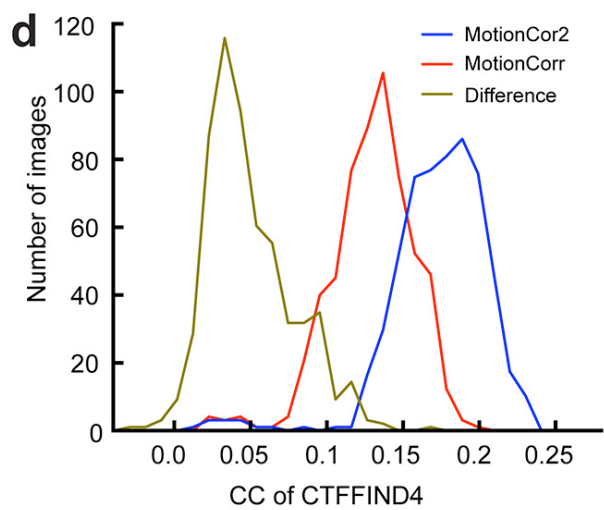
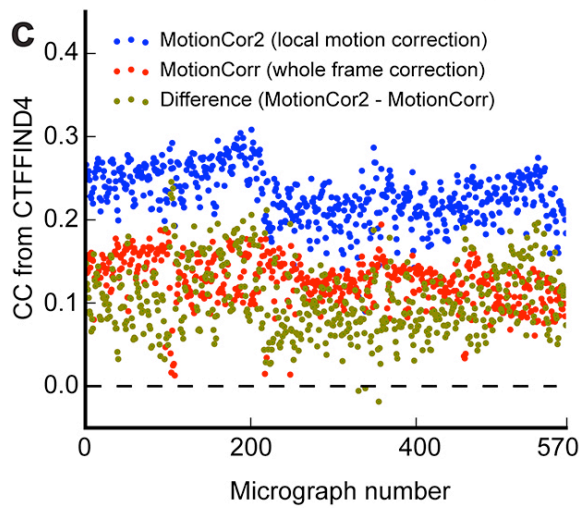
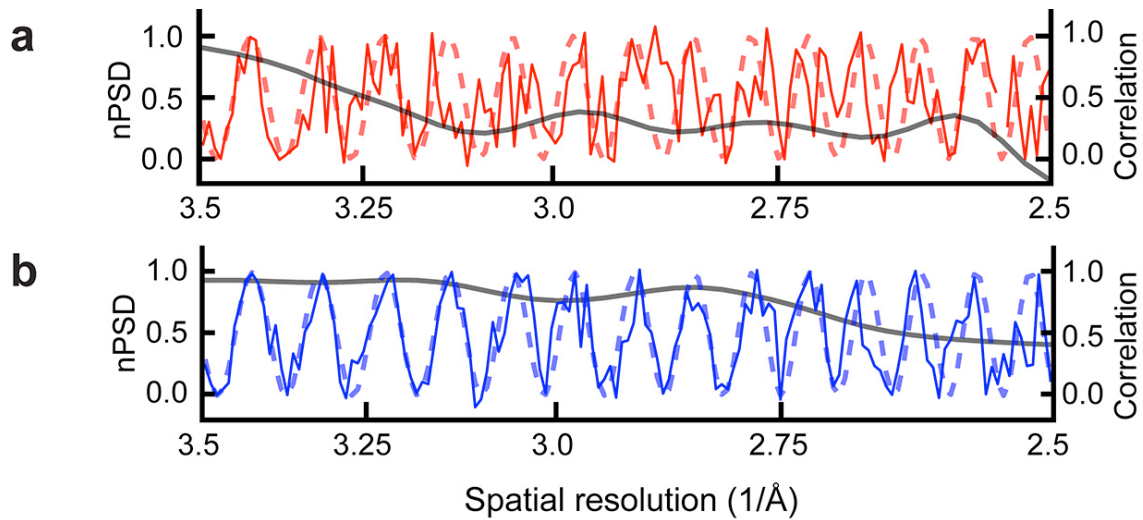


Supplementary Figure 3

Influence of dose weighting on the Fourier power spectrum

(a) and (b) Fourier power transform calculated from dose-weighted (a) and un-weighted (b) image after motion correction. (c) The rotation averages

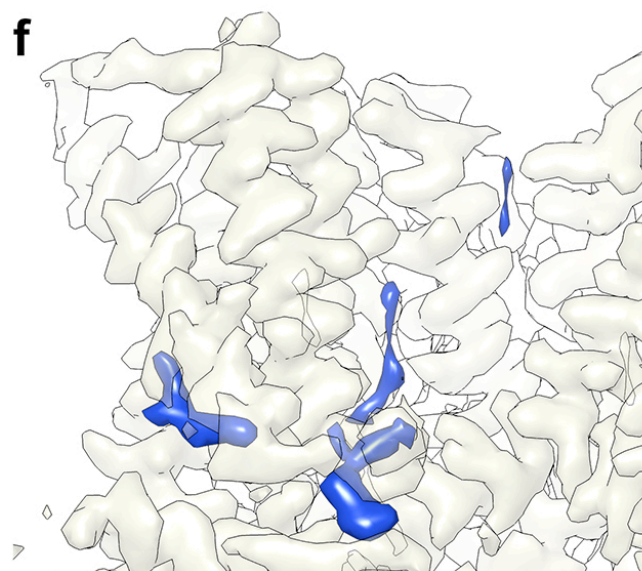
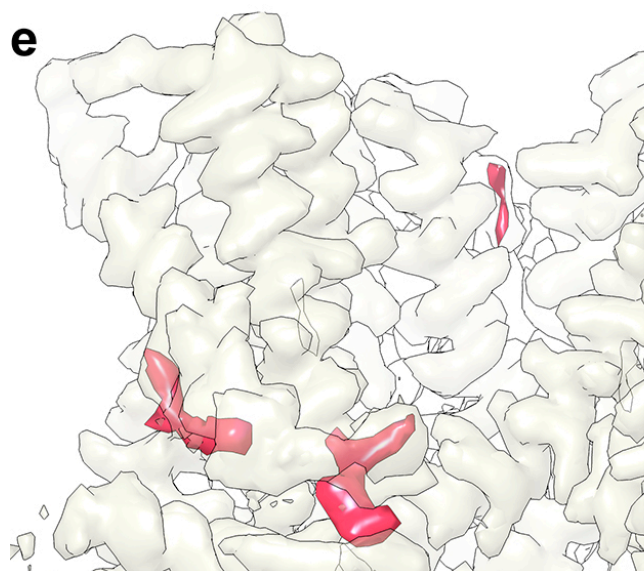
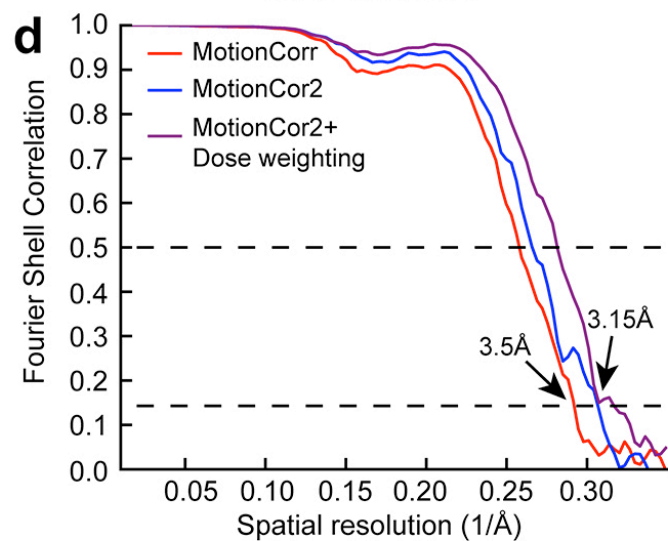
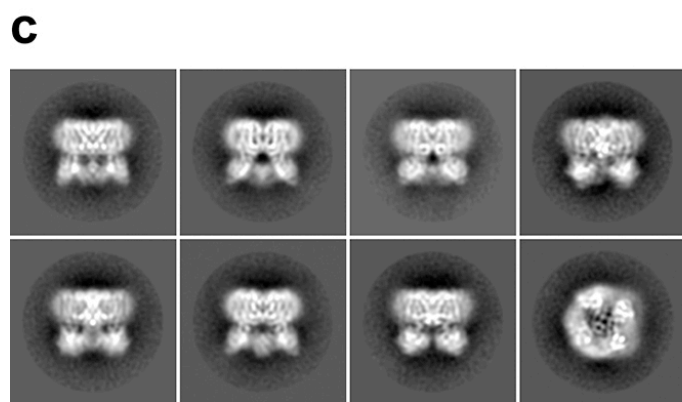
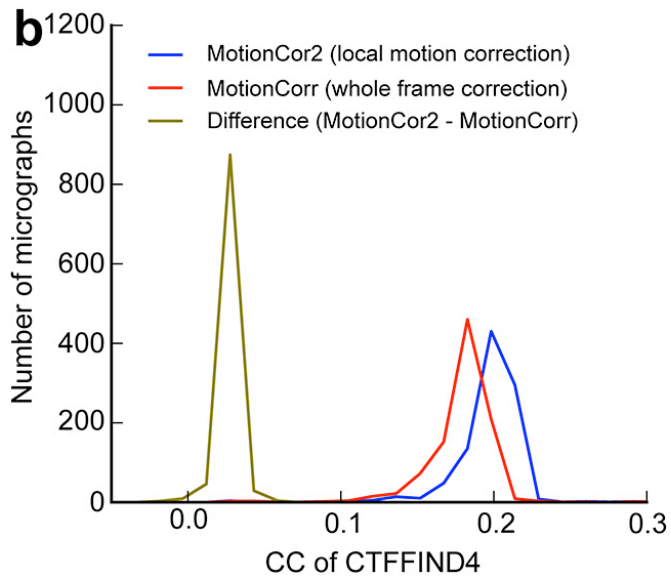
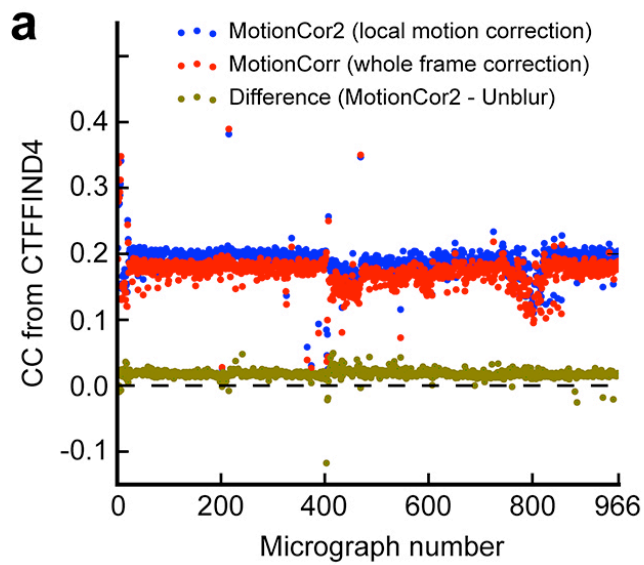
of dose-weighted (blue) and un-weighted (red) Fourier power spectra shown in (a) and (b).



Supplementary Figure 4

Comparison of motion corrections by MotionCor2 and MotionCorr on 3D reconstruction of archaeal 20S proteasome.

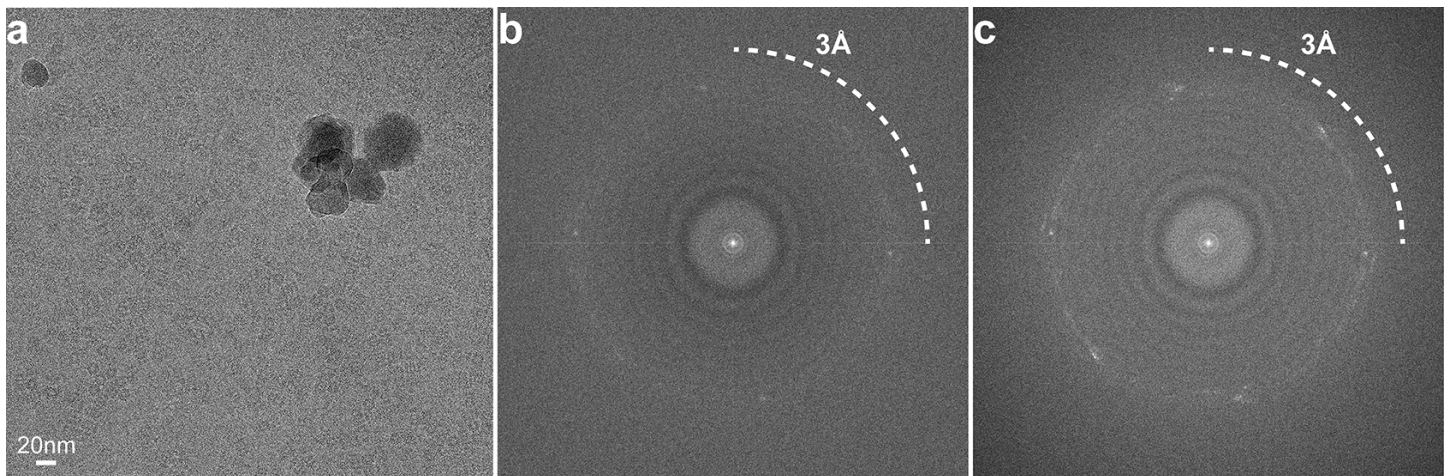
(a) Rotationally averaged Fourier power spectrum of image after motion correction by MotionCorr (solid line) and fitted contrast transfer function (dashed line). (b) Rotationally averaged Fourier power spectrum of image after motion correction by MotionCor2 (solid line) and fitted CTF (dashed line). Solid black line in both (a) and (b) indicate cross correlation coefficient between the rotationally averaged Fourier power spectrum of image and fitted CTF. (c) The cross correlation coefficient output from CTFFIND4 was used to assess the Thon ring quality within the resolution range of 10 ~ 5 Å. Blue and red dots represent cross correlation coefficient of every micrograph after motion correction using MotionCor2 (blue) and MotionCorr (red). Brown dots represent the differences between the two. (d) Histogram of cross correlation coefficients between calculated and simulated Fourier power spectrum of MotionCorr corrected image (red) and MotionCor2 corrected image (blue). The difference, which shows the amount of improvement, is shown in brown. (e) Comparison of motion corrections by MotionCor2 and Unblur. The cross correlation coefficient output from CTFFIND4 calculated within the resolution range of 10 ~ 5 Å. Blue and red dots represent cross correlation coefficient of every micrograph after motion correction using MotionCor2 (blue) and Unblur (red). Brown dots represent the differences between the two. (f) Histogram of cross correlation coefficients between calculated and simulated Fourier power spectrum of MotionCor2 corrected image (blue) and Unblur corrected image (red). The difference, which shows the amount of improvement, is shown in brown. (g) Selected 2D class averages of 20S proteasome calculated from images after motion correction by MotionCor2.



Supplementary Figure 5

Comparison of motion correction by MotionCor2 and MotionCorr on 3D reconstruction of rat TRPV1 ion channel.

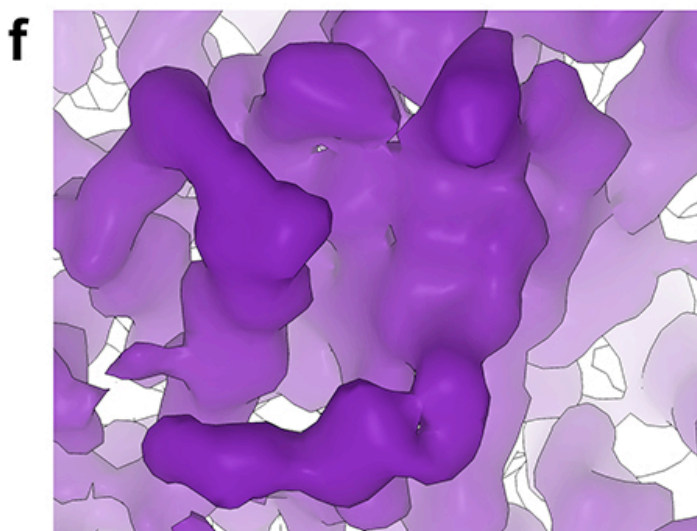
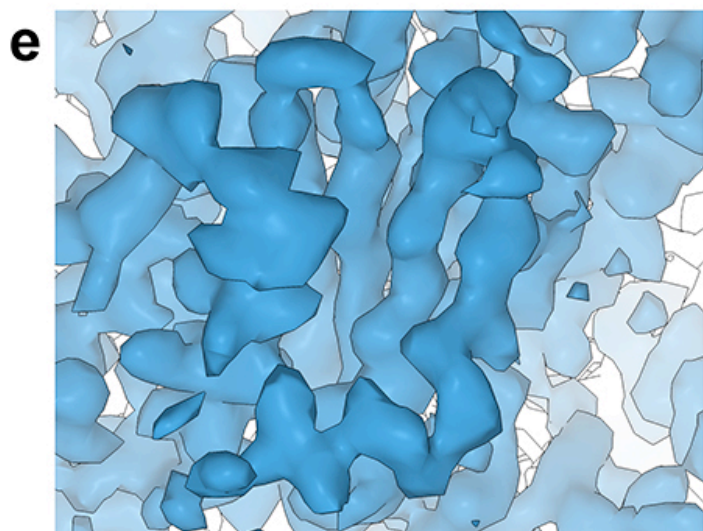
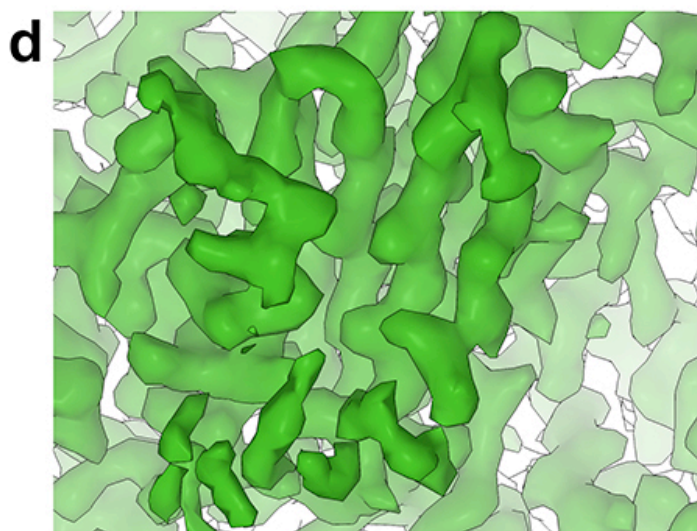
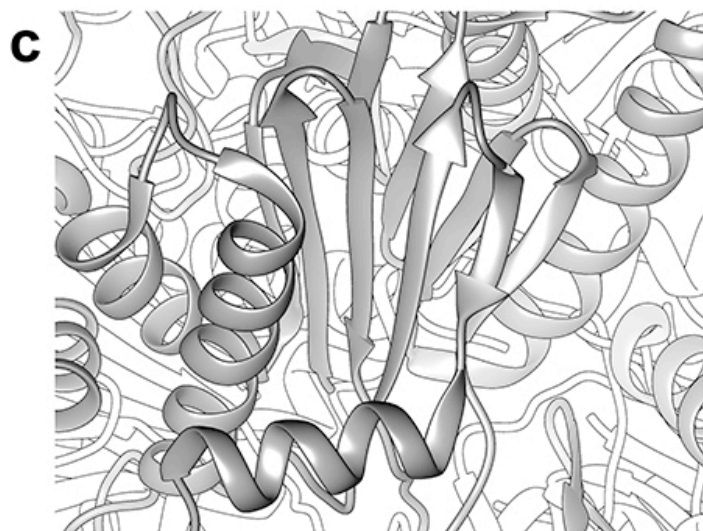
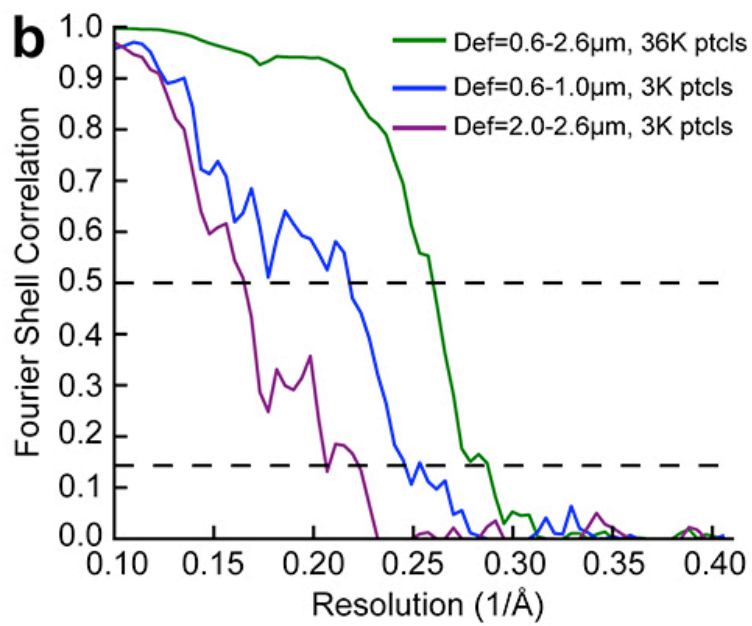
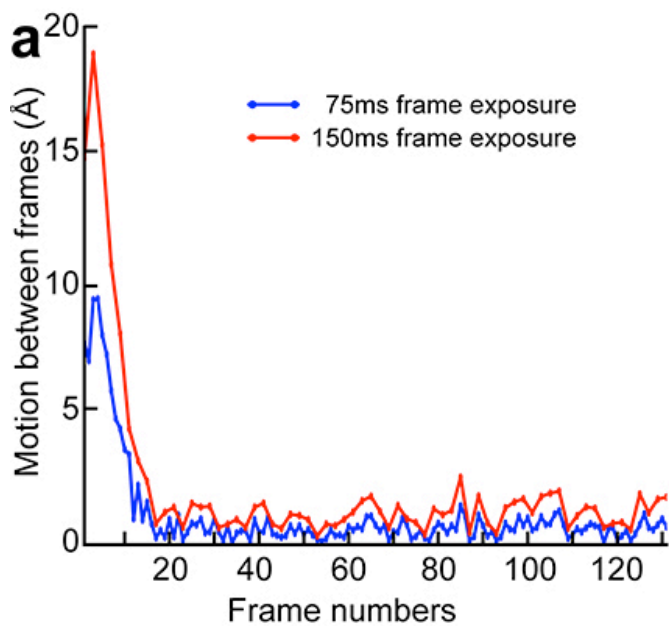
A published dataset of frozen hydrated rat TRPV1 ion channel was reprocessed using MotionCor2. (a) Cross Correlation Coefficient of CTFFIND4⁹ from the correction by both MotionCorr (red) and MotionCor2 (blue). (b) Histogram of cross correlation coefficients determined by using CTFFIND4 using image corrected by MotionCor2 (blue) and MotionCorr (red). The difference, which shows the amount of improvement, is shown in brown. (c) Representative 2D class averages of frozen hydrated TRPV1 particles. (d) FSC curves of 3D reconstructions determined from the same dataset after motion correction by MotionCorr (red), MotionCor2 (blue) and MotionCor2 with dose weighting (brown). (e) A representative view of the TRPV1 ion channel generated from previously published density map⁵. (f) The view of the same region of TRPV1 density map determined after re-process motion correction using MotionCor2. Both maps (e and f) are shown at the same normalized density level, $\sigma = 6$.



Supplementary Figure 6

Motion correction of low defocused image.

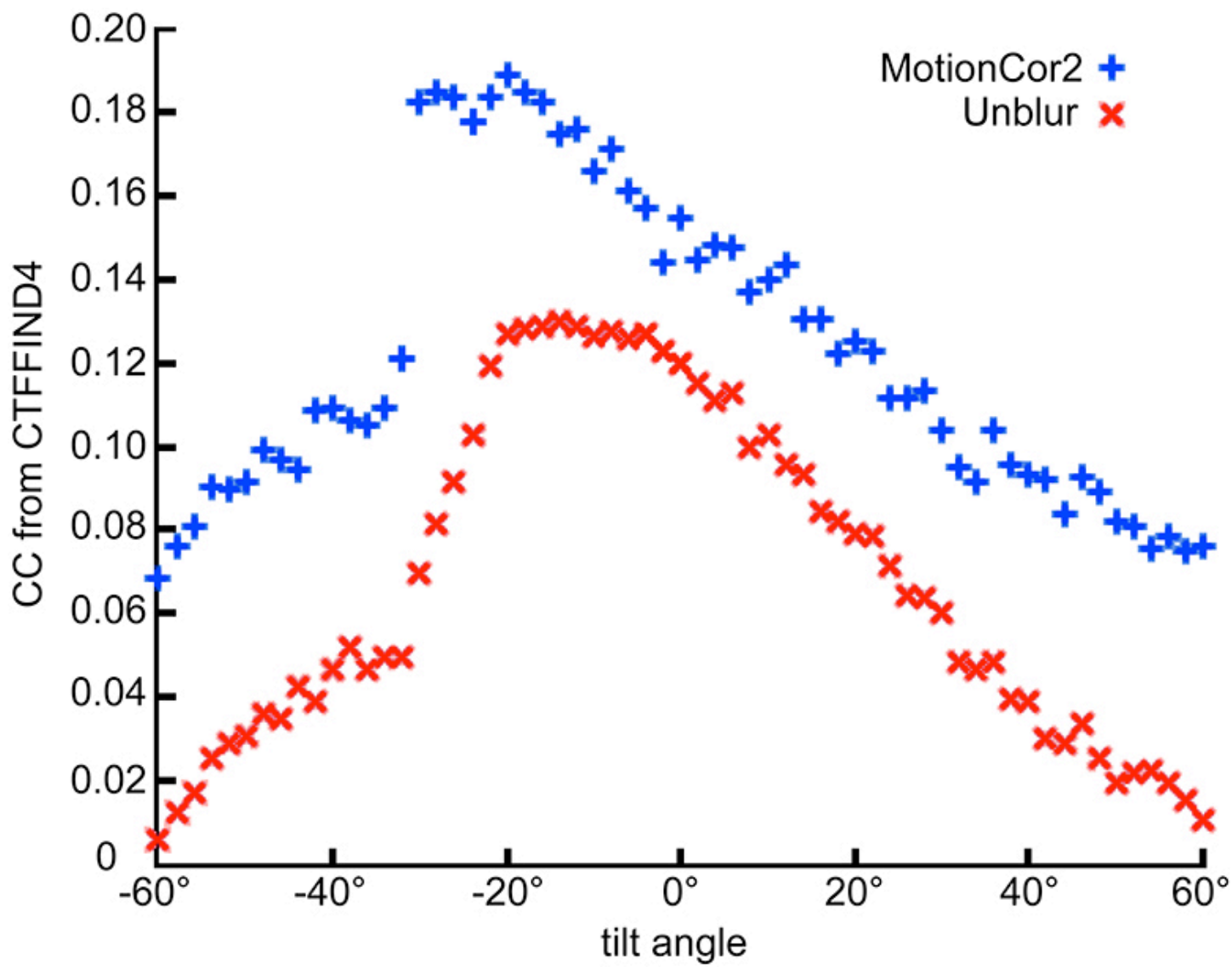
(a) A micrograph of frozen hydrated archaeal 20S proteasome recorded with a defocus of $0.4\ \mu\text{m}$ at 200kV on a TF20. (b) and (c) Fourier power spectra are shown from the same image before motion correction (b), and after motion correction with MotionCor2 (c). The motion correction of the same micrograph by MotionCorr failed.



Supplementary Figure 7

3D reconstruction of archaeal 20S proteasome from a TF20 200kV electron microscope.

(a) A plot of motion between neighboring frames when frame exposure was set to 0.075 second (blue) and 0.15 second (red, by average two adjacent sub-frames to produce a sub-frame that is equivalent to the 0.15 second frame exposure time). (b) Red: FSC curve of an archaeal 20S proteasome 3D reconstruction determined from a dataset of ~36,000 particles collected with a TF20 electron microscope operated at 200kV acceleration voltage. The defocus range was set between 0.6 μ m and 2.6 μ m. Blue: FSC curve of a 3D reconstruction using a subset of 3,000 particles with low defocus (0.6 μ m to 1 μ m). Purple: FSC curve of a 3D reconstruction of using another subset of 3,000 particles high defocus (2.0 μ m to 2.6 μ m). (c) Ribbon diagram of a part of archaeal 20S proteasome. (d) Same region of the 3D reconstruction determined from the entire dataset, corresponding to red FSC curve in (a). (e) Same region of the 3D reconstruction determined from the subset of 3,000 particles with only low defocused particles (blue FSC curve). (f) Same region of the 3D reconstruction determined from a subset of 3,000 particles with only high defocused particles (purple FSC curve). All maps (d, e and f) are shown at the same normalized density level, $\sigma = 4$.



Supplementary Figure 8

Comparison of Unblur and MotionCor2 on a tomographic tilt series.

This tomographic tilt series was collected between $[-60^\circ, 60^\circ]$ at every 2° from a frozen hydrated specimen of *Drosophila* centriole. The data collection was performed in two branches that started from -30° to 60° and then from -32° to -60° and at the magnification that gives rise to the pixel size of 4.08\AA . At each angular step a movie stack of 20 sub-frames was collected at dose rate of $5e^-/\text{pixel}/\text{s}$ with frame exposure of 0.25s . Under this setup the per-frame dose is $0.075e^-/\text{\AA}^2$ and the total dose of the entire tilt series is $92e^-/\text{\AA}^2$.

MotionCor2 - anisotropic correction of beam-induced motion for improved cryo-electron microscopy

Shawn Q. Zheng^{1,2}, Eugene Palovcak¹, Jean-Paul Armache¹, Yifan Cheng^{1,2} and David A. Agard^{1,2}

¹Department of Biochemistry and Biophysics and

²Howard Hughes Medical Institute, University of California, San Francisco, CA 94158

Supplementary Table 1: Comparison of atomic models of archaeal 20S proteasome (*T. acidophilum*) refined against the 3D density maps determined previously⁴ and in this work.

	Previous work	Current work
Total number of particles:	129,008	187,011
Map resolution (FSC=0.143)	3.3Å	2.58Å
Molprobity score	2.55	1.81
Ramachandran outliers	1.9%	0.00%
Ramachandran favored	90.10%	96.2%
Rotamer outliers	9.30%	0.6%
C-beta outliers	574	0
Clash-score	5.11	10.79
Real-space cross correlation	0.75	0.76
EMRinger score	3.08	3.95

SUPPLEMENTARY METHODS

Background Sample illumination with the high-energy electron beam breaks bonds, releases radiolysis products and builds up charge within the thin frozen hydrated biological samples during image recording. The result is a combination of physical and optical distortions that can significantly deteriorate sample high-resolution information through image blurring. This has been one of the major factors limiting the achievable resolution of single particle cryo-EM. The concept of recording the image as movie to correct sample motion was proposed long ago^{10,11}, but only became practical once direct electron detection cameras became available¹². The fast sample motion can be measured by tracking the movement of images captured in a series of snapshots, either as the whole frame or as individual particles. Image motion can then be corrected by registering identical features in the sub-frames to each other, followed by summing the registered sub-frames to produce a motion-corrected image. While in principle tracking should be straightforward, the practical challenge is the extremely low signal-to-noise ratio (SNR) in each individual sub-frame. Except for very large particles, accurate motion measurement generally requires correlating the motion between sub-frames over large areas. That said, even sub-optimal motion correction can significantly restore high-resolution signals and improve the resolution of final 3D reconstructions.

We previously developed an algorithm to correct uniform image motion. It uses redundant measurements of image shifts between all sub-frames to derive a least squares estimate of relative motions between neighboring sub-frames. We implemented this algorithm in the program MotionCorr that provided an efficient correction of image motions with sufficient accuracy⁴ to enable the determination of numerous near atomic resolution 3D reconstructions^{5,13}. Around the same time or soon afterwards, a number of different strategies were devised that either assume particles located nearby have similar motions or assume uniform motion of the entire frame or patches of the frame. Programs based on the former assumption include RELION that provides a movie-processing mode but uses a 3D reconstruction to track particle motions^{6,14}, Xmipp that implemented an Optic Flow algorithm¹⁵, and *alignparts_lmbfgs* that implemented a regularized Fourier Space optimization algorithm to track neighboring particles⁷. Programs based on tracking the full frame or parts of the frame include MotionCorr⁴ and iterative whole frame alignment procedures such as Unblur^{8,16} or those used in electron tomography¹⁶. All of these

algorithms have demonstrated the ability to recover high-resolution signals to varying degrees and have improved the resolution of the resultant 3D reconstructions.

Ideally, single particle cryo-EM images should be acquired with the smallest possible defocus to enhance high-resolution information, and with the shortest possible sub-frame exposure times to reduce motion trapped within individual sub-frames, in particular for the first few frames where the sample has the least radiation damage but moves most rapidly¹⁷. Additionally, motion detection should be done on the smallest possible local area to best capture the anisotropic motion. Unfortunately, minimizing defocus, time or area significantly reduces the SNR available for measurement, ultimately leading to incorrect motion estimates. For example, previous experiments with MotionCorr revealed that subdividing the images into areas smaller than $\sim 2000 \times \sim 2000$ pixels or going to sub-frame integration times of less than 100 milliseconds worsened resolution due to increased errors in motion tracking.

In this work we experimentally validated the doming behavior that was proposed to describe beam-induced sample motion³ and discovered that the initial motion begins at surprisingly low doses. Our observations supported a model in which the sample is smoothly deformed throughout the exposure. The extreme sensitivity to dose emphasizes the need to acquire at very short sub-frame durations, further raising the demands for performance at low SNRs.

To meet the compounding demands for locality and low SNR, it was not possible to accurately determine local shifts by comparing noisy sub regions of each frame to one another as done within MotionCorr. Instead, we compare local image regions to an approximated sum for each image subregion, and then iterate translation determination to improve the sum. The resultant shifts are further constrained both spatially and temporally as informed by the physical behavior of the sample. This is accomplished by deriving the final shifts from a fit of a time-varying two-dimensional (2D) polynomial function to the local motions derived from the different patches of the image. Each image sub-frame is subsequently remapped using this smooth distribution of beam induced motions at each individual pixel and summed with or without radiation damage weighting. This algorithm has been implemented as MotionCor2, a parallel computing program running on Linux platform equipped with multiple GPUs. Our tests have shown that this program is very robust, and sufficiently accurate at correcting local motions

so that the very time-consuming and computationally-intensive particle polishing procedure in RELION can be skipped.

Modeling of beam-induced motion Early studies suggested that beam induced motion can be described as specimen doming³. This model predicts a significant sample motion perpendicular to the sample plane (z-motion, Figure 1a), which, when the sample is tilted, should be projected onto the image plane as a motion perpendicular to the tilt axis (Supplementary Figure 1a). We explored this by directly measuring motions within movie stacks recorded from specimens tilted at various angles. We collected a number of dose fractionated tomographic tilt series from various frozen hydrated biological samples. By using rather low magnifications (4.3Å/pixel), we could accurately determine full frame shifts with doses down to 0.05 e⁻/Å² using MotionCorr⁴. Representative traces of uniform whole frame motion of the same specimen at three tilt angles $\alpha = -30^\circ$, 0° , and 60° clearly show that the measured motion is much larger at higher tilt angles and mainly perpendicular to the tilt axis (Supplementary Figure 1b), supporting the dome model. Consistent with an earlier study³, the motion within the very first 0.2 e⁻/Å² is substantial. The beam-induced shifts observed during normal single particle cryo-EM must be a combination of sample drift, residual tilt, and non-idealities in the doming such that local motions are not purely along Z.

Motion correction using polynomial constraints Whole frame based motion correction algorithms, such as^{4,8,16}, correct only the uniform motion across the entire image frame, leaving the residual non-uniform local motion uncorrected. A common strategy for correcting local motions is to spatially partition a movie stack into multiple stacks of patches where the correction is performed individually⁴. Ideally, the smaller the patch size, the better the locality of motion correction. In practice, there is a limit on patch size, below which the motion cannot be corrected with sufficient accuracy due to insufficient SNR. Furthermore, dividing the whole image into a limited number of patches and correcting motion within each patch independently causes edge artifacts at the boundaries of patches in the corrected image, the so-called checkerboard artifact. This not only creates additional challenges in single particle work, but obviates utilization for tomography.

To avoid these problems, we chose to combine iterative patch-based motion measurement with restraints derived from the physical behavior to obtain a function that describes the motion of each pixel across the entire image. Since a dome can be geometrically approximated by a quadratic surface, we chose to fit the locally measured x and y shifts from all patches independently to a time-varying polynomial function that is quadratic in the xy plane and cubic with respect to exposure time t .

$$\begin{aligned}
S(x, y, t) = & c_0t + c_1t^2 + c_2t^3 + c_3xt + c_4xt^2 + c_5xt^3 + c_6x^2t + c_7x^2t^2 + c_8x^2t^3 \\
& + c_9yt + c_{10}yt^2 + c_{11}yt^3 + c_{12}y^2t + c_{13}y^2t^2 + c_{14}y^2t^3 \\
& + c_{15}xyt + c_{16}xyt^2 + c_{17}xyt^3
\end{aligned} \tag{1}$$

where, S is the shift at the integer pixel coordinate x and y . At any given time point τ , the above equation becomes a standard quadratic function:

$$S(x, y)_{t=\tau} = a_0 + a_1x + a_2x^2 + a_3y + a_4y^2 + a_5xy \tag{2}$$

that approximates a domed surface. Since Eq. (1) is high-order differentiable over x , y and t , a smooth estimate of the shift can be obtained at each pixel of each sub-frame and between consecutive sub-frames. Not only does this strategy eliminate checkerboard artifacts, but importantly it also provides a more thorough correction of image motion. It effectively dampens the fluctuations of local measurements influenced by image noise.

Motion correction algorithm Image motion can be decomposed into two components, the uniform global motion and the non-uniform local motion that is described as the projection of the doming motion. The motion correction is divided into two steps, global motion correction followed by local motion correction.

The global motion is corrected iteratively by calculating cross correlation of each sub-frame against the sum of all other sub-frames based upon the alignment of the previous iteration. To avoid self-correlation, the sub-frame being measured is excluded from the sum. The measured shifts, i.e., the residual errors of the previous iteration of alignment, are corrected by shifting the phase in the Fourier transforms of sub-frames. The iteration stops when the residual errors are below a specified tolerance (typically, 0.5 super-resolution pixel) or a specified maximum number of iterations have been reached.

The global-motion-corrected stack is partitioned into non-overlapping patches. The same alignment procedure is applied to each stack of patches. The local motion is therefore measured at a series of discrete spatial locations represented by the patch centers. The measured local shifts of each patch for each sub-frame are fit to a polynomial function given in Eq. (1). Least squares fitting is used to determine the unknown coefficients c_i . The measured x and y shifts are fit separately to Eq. (1), resulting two polynomial functions, $S_x(x,y,t)$ and $S_y(x,y,t)$, corresponding respectively to the x and y components of local motion. The vector function (S_x, S_y) provides a complete and smooth description of the local motion field, allowing smooth correction of the local motion at the single pixel level without causing edge effects.

Since the calculated shifts are non-integer and varying from pixel to pixel, the correction of pixel shifts must be performed in real space and we used bilinear interpolation. To minimize the attenuation of high-resolution signals due to interpolation, the interpolation is performed on super-resolution pixels. The final image is obtained by cropping the corrected sum in Fourier domain to the user specified resolution.

Bad Pixel correction Despite the considerable efforts of the camera manufacturers, there are often image defects that remain uncorrected by their calibration methods. In the case of our Gatan K2 cameras, we observe a variety of imperfections including hot pixels, cold pixels and clusters of 6×6 super-resolution pixels having distorted intensities (Supplementary Fig 2a). As the intensity distribution is similar to a 2D truncated sinc function (or an inverted version), these are likely the consequence of a single pixel defect perturbed by Gatan's electron counting/super-resolution algorithm. It was also observed that the locations of the defects often vary from exposure to exposure (Supplementary Fig 2b and c). The longer the exposure, the more clusters are present. To minimize possible deleterious effects of these systematic image errors on the low SNR image registration, we first detect these defects based on the image statistics within each stack. Defects are identified first from the sum of gain-corrected sub-frames without motion correction, with the clusters located by cross-correlating a template 2D truncated sinc function. Once identified, these defects are corrected for each sub-frame using values of good pixels randomly selected from the neighborhood. Only the defect-corrected stacks are used for subsequent motion correction.

Motion correction program MotionCor2 We first correct the uniform global motion by iteratively aligning each sub-frame against a reference that is the motion corrected sum of all other sub-frames based upon the translational alignment obtained in the previous cycle. Excluding the sub-frame being aligned from the reference prevents a strong auto-correlation peak at the origin that may influence the determination of the real cross correlation peak. The measured sub-frame shifts at each iteration are the residual errors from the last iteration, and the alignment procedure is terminated when the maximum residual error is below a specified tolerance. The measured global shifts are then corrected by phase shifting in the Fourier domain to yield a global-motion-corrected stack. The image stack is then divided into a grid of patches, and the same iterative alignment procedure is performed on each patch. This provides the local motion at a series of discrete locations for each time point within the exposure. These patch based local motions are then fit to polynomial functions defined in Eq. (1). As a result, x and y shifts can be calculated according to the fit functions at each pixel (m, n) in each sub-frame. Since the calculated shifts at individual pixels are typically non-integer, the corresponding correction requires interpolation in real space. To minimize the attenuation of high-resolution signals due to interpolation, bilinear interpolation is performed on super-resolution pixels. The final image is obtained by cropping the corrected sum in Fourier domain to the user specified resolution. In practice the whole stack is typically divided into 5×5 partitions, an empirical choice that provides a good balance between precision, efficiency, and SNR. However, for the higher contrast Volta phase plate images, good results have been obtained using 9×9 patches¹⁸.

In order to maximize the throughput of motion correction, parallel computation was implemented in three levels in MotionCor2 on a Linux platform equipped with multiple GPUs.

Integration of dose weighting Dose weighting each motion-corrected sub-frame according to expectations from analysis of radiation damage allows cryo-EM images to be recorded with significantly higher total electron doses. This improves low-resolution image contrast without sacrificing high-resolution SNR⁸. We have integrated this weighting scheme into MotionCor2 to streamline all the necessary preprocessing steps prior to the normal cryo-EM processing pipeline.

For biological samples embedded in vitreous ice, radiation damage dampens high-resolution signals that are useful for high-resolution structure determination, hence better results should be obtainable via proper weighting. While the dose-weighting scheme appropriately

down-weights the high-resolution biological information from the high dose data within the summed image stack, it also unnecessarily attenuates the high-resolution Thon ring signals used for CTF determination. To avoid this undesirable side effect, MotionCor2 generates both a dose-weighted summed image for cryo-EM reconstruction and an un-weighted summed image for CTF estimation (Supplementary Figure 3).

Parallel implementation Our constrained, patch-based motion correction algorithm was implemented in a GPU (Graphics Processing Unit) accelerated program named MotionCor2 that runs on a Linux platform equipped with one or more GPUs. High-throughput motion correction was achieved by means of parallel computation implemented in three levels. At the lowest level, pixel-wise computations, such as correction of gain reference to each pixel of a sub-frame, are implemented in various CUDA kernels. At the second level, independent operations such as Fourier transform of sub-frames and alignment of individual patches are distributed to all participating GPUs. The last level of parallelization is the batch processing of multiple movie stacks. The disk IO operations, i.e. loading a new stack and saving a corrected image, are performed in parallel with the intensive computation involved in correcting a loaded stack. This strategy minimizes the idle time of GPUs in waiting for a new stack. We assessed the efficiency of MotionCor2 configured for 5×5 patch-based motion correction on a Linux platform equipped with 4 Tesla K10 GPUs (NVIDIA). The input movie stacks contain 30 sub-frames of 7676×7420 pixels of 8-bit pixel depth. The corrected images were truncated to 3838×3710 pixels in Fourier space before they are saved to disk. For a single stack it took a total of ~62 seconds of which ~42 seconds were spent on computation with the remaining 20 seconds for disk operations. For a batch correction of 10 such stacks, it took ~444 seconds in total. On average the correction of each stack took only 44 seconds. The majority of disk IO time is shadowed by the computational time.

MotionCor2 improves accuracy of motion correction We tested the performance of MotionCor2 using two previously acquired single particle cryo-EM datasets: the archaeal 20S proteasome and the rat TRPV1 ion channel. Both were previously processed using whole frame based MotionCorr and produced near atomic resolution structures^{4,5}. These two datasets were collected using the same microscope settings and similar electron dose rates on the camera. As these were

older data sets the total dose was limited: the 20S stacks contained 25 sub-frames (total dose of ~ 34 electrons per \AA^2) whereas the TRPV1 stacks had 30 sub-frames (total dose of ~ 41 electrons per \AA^2). We re-processed motion correction for these two datasets using MotionCor2, configured to run on 5×5 patches.

The trajectory of the full frame (global) motion determined by the original MotionCorr program is shown in Supplementary Figure 1c. A very similar global motion trace was determined using the MotionCor2 iterative strategy (long trace, Supplementary Figure 1d). The measured local motions for each patch are more complex (traces in each panel). In this representative image stack the smoothed trajectories (Figure 1b) are qualitatively similar to the raw traces (Supplementary Figure 1d) indicating that the added robustness of applying physical constraints does not come at the cost of capturing local motion. The new algorithm significantly improves Thon ring signals at high-resolution and leads to a better correlation with the simulated contrast transfer function (CTF) (Supplementary Figure 3c, Supplementary Figures 4a,b). Over the entire data set, almost all images showed a significant improvement in recovery of high-resolution Thon ring signals (Supplementary Figure 4). Performing a similar analysis with Unblur, a very popular current generation iterative whole frame motion correction strategy, shows that MotionCor2 provides a further improvement in Thon ring resolution (Supplementary Figure 5).

To better understand the impact of the improved motion correction, we re-determined the 3D reconstructions of the archaeal 20S proteasome from images that were processed by MotionCorr, Unblur and MotionCor2, followed by refinement and reconstruction with the maximum-likelihood based refinement and reconstruction program, RELION (Figure 1c).

In the 3D reconstruction of T20S after MotionCor2 and dose weighting (green FSC in Figure 1c) most backbone carbonyls are now clearly visible, as well as the precise rotameric states of amino acid side chains (Figure 1d - f). Such fine structural detail would clearly facilitate more accurate model building. Indeed, after real-space refinement of our previous atomic model into the new map, we have substantially improved model validation statistics such as the EMRinger score, the real-space cross correlation, and various model geometry scores given by MolProbity (Supplementary Table 1).

Similarly, we re-processed our previous raw micrographs from the TRPV1 ion channel⁵ and re-determined its 3D reconstruction. MotionCor2 improves the nominal resolution from

3.5Å to 3.1(5)Å (Supplementary Figure 5). Such improvements are particularly obvious in some trans-membrane regions, where extra densities associated with the TRPV1 protein are now seen to have well-defined features that can only now be interpreted as lipid molecules (Supplementary Figure 5e and f). Thus, for both the 20S proteasome and TRPV1 datasets, MotionCor2 produced noticeable resolution improvements.

We also tested the performance of MotionCor2 with a new dataset of archaeal 20S proteasomes collected from our TF20 microscope operated at 200kV equipped with a Gatan K2 camera. The sub-frame exposure time was set to 75ms, approximately one third of the frame exposure time (200ms) we typically use. The motion captured within each 75ms sub-frame, particularly in the early sub-frames, is significantly less than that captured within a sub-frame with a longer frame exposure time (Supplementary Figure 7a). About half of these images were recorded with defocus set between 0.4µm to 1.0µm. A representative image collected with 0.5µm defocus is shown in Supplementary Figure 6 and MotionCor2 restored its Thon rings to close to 3Å resolution. We picked a small set of 36,000 particles from motion corrected images and determined a 3.4Å reconstruction of the 20S proteasome (Supplementary Figure 7b). Interestingly, a reconstruction of 4.1Å resolution could be determined using only the 3,297 particles having defocus values less than 1 µm (Supplementary Figure 7b, c and d). In contrast, using the same number of particles, but higher defocus (2 ~ 2.5 µm), the resolution of the 3D reconstruction is considerably worse, 4.9Å (Supplementary Figure 7b and e). This emphasizes the importance of being able to accurately motion correct very low defocus images.

Tomography data A major advantage of the algorithm implemented in MotionCor2 is to correct image motion captured in electron tomography tilted series, where the approach of tracking individual particles is not applicable. Thus, the algorithm implemented in MotionCor2 is currently the only way to deal with local image motion captured in movie stacks of tilted specimens. As shown in Supplementary Figure 8 using a representative tilt series of frozen hydrated *Drosophila* centrioles, MotionCor2 restores image Thon ring from tilted specimen significantly better than the full frame correction made by Unblur.

Further discussion Improving both accuracy and efficiency of motion correction has been an extensive topic of research since direct electron detection cameras and dose-fractionated imaging

techniques were first introduced to cryo-EM^{4,6,7,14-16}. Compared with algorithms that only correct the global motion, the new algorithm implemented in MotionCor2 has shown consistent and significant improvements while maintaining superb robustness. Because image motions are corrected at the pixel level prior to any further image processing, it is practical to take advantage of the benefits of dose-weighting without the computationally intensive steps that track and correct motion of individual particles implemented in RELION¹⁴. Importantly, our full frame motion correction strategy is also applicable to tilted data collected by cryo tomography. Furthermore, by outputting both unweighted and dose-weighted frame averages, CTF determination is optimized. Particles boxed from the weighted images have enhanced low-resolution contrast and high-resolution SNR, facilitating more accurate particle alignment and classification, thereby producing 3D reconstructions having improved resolutions.

Our tests have shown that the new algorithm also works well with dose-fractionated image stacks recorded with defocus and frame exposure times that are significantly lower and shorter than commonly used (Supplementary Figure 7). This could potentially impact the collection of single particle cryo-EM datasets when aiming for the highest resolution possible. For atomic-resolution single particle cryo-EM reconstruction, there are many benefits to recording images at low defocus, preferably less than 1 μm ¹⁷. Provided that image motion can be corrected accurately, dose weighting allows images to be recorded with sufficiently high total dose to enable processing while realizing the benefits of low defocus to obtain high resolution. Another practical issue is that motion is most rapid and most anisotropic in the first few sub-frames, resulting in substantial deterioration of high-resolution information within the first $\sim 3 \text{ e}^-/\text{\AA}^2$ of dose. While these have typically been excluded or substantially down weighted, we have shown that recording images with shorter sub-frame exposure times can reduce such deterioration, again, provided that the local image motion captured in the short sub-frames can be corrected (Supplementary Figure 7a). Our tests have shown that MotionCor2 provides stable correction of images recorded at both low defocus values and short frame exposure times (Figure 4). Such robust motion correction enabled determination of a 4 \AA resolution reconstruction using only $\sim 3,000$ particles from low-defocused images with short frame exposure times recorded on a TF20 microscope. This dramatically enhances the utility of such an instrument for routine single particle cryoEM.

Our tests (Figure 1c) suggest that after motion correction by MotionCor2, it is possible to skip the particle tracking part of the RELION polishing procedure, which is by far the most computational intensive and time consuming step in the refinement of a 3D reconstruction. However, these same experiments suggest that there can be some further benefit from the individual frame weighting component of the RELION polishing procedure. It is perhaps not too surprising that the dose weighting scheme, which was originally implemented in Unblur, may not be as accurate as RELION per-frame B-factor weighting, which is determined individually for each dataset. This may also suggest that factors other than pure radiation damage (the model in Unblur) may be affecting the mid and high dose data.

The algorithm implemented in MotionCor2 shares a conceptual link with the unbending procedure used for processing images of two-dimensional crystals¹⁹, which involves local interpolations with a long range constraint imposed by the crystallinity. Similarly, the x,y, and time-dependent polynomial function used in MotionCor2, also constrains the interpolation, minimizing movements of individual pixels to reduce noise. For typical images of both T20S proteasome and TRPV1, the largest difference of measured shift within a 4,320 x 4,320 Å specimen area is less than ~12Å. Under this estimation, a particle of ~300Å, which is about the size of an eukaryotic ribosome or spliceosome particle, could be distorted by up to ~0.8Å. For very large objects, tracking individual particles might be a better strategy. However, for smaller objects, the accuracy, robustness and computational convenience of MotionCor2 provides a significant advantage.

Cryo-EM data acquisition Raw micrographs of archaeal 20S proteasome (*T. acidophilum*) and TRPV1 ion channel were acquired in previous studies, as described^{4,5}. We also collected a new dataset of archaeal 20S proteasome using Tecnai TF20 electron microscope (FEI Company) equipped with a field emission electron source operating at 200kV acceleration voltage and K2 Summit camera operating in super-resolution mode (Gatan). Frozen hydrated 20S proteasome grids were prepared as previously described. A total of 141 micrographs were collected at a magnification of 29,000X has a physical pixel size of 1.234 Å. The dose rate was set to ~8 e⁻/physical pixel/second at camera level. Frame exposure time was set to 0.075 second, and total exposure at 10 seconds, corresponding to a total accumulative electron dose of ~65 e⁻/Å² on specimen. The defocus range was set to 0.5 ~ 2.5µm.

Image processing Images were subject to motion correction using MotionCorr, Unblur and MotionCor2. After motion correction, images were 2x binned by Fourier cropping implemented in both programs. Un-weighted sums were used for CTF determination using CTFFIND4. Weighted sums were used for automated particle picking using RELION and subsequent image processing. The first sub-frame was excluded in both cases. In MotionCor2, we used per-frame dose of $1.36 \text{ e}^-/\text{\AA}^2$ (20S proteasome), $1.30 \text{ e}^-/\text{\AA}^2$ (TRPV1), and $0.58 \text{ e}^-/\text{\AA}^2$ for the new proteasome dataset collected from TF20 microscope for calculating the critical exposure curves for dose-weighting.

Normalized elliptically-averaged Fourier spectrum of un-weighted sum of motion-corrected sub-frames was used to CTF determination using CTFFIND4⁹. In each dataset, we used 2D class averages calculated from a small number of manually picked particles as references for automatic particle picking using RELION²⁰. 3D reconstructions of each dataset were also calculated and refined following gold-standard refinement procedure implemented in RELION.

For 20S proteasome dataset collected at 300kV, a total of 221,623 particles were auto-picked from 570 micrographs corrected with MotionCor2. After reference-free 2D classification, particles from classes that did not show high-resolution features were excluded, leaving a final set of 187,011 particles. The coordinates of this particle set were used to extract the equivalent particle images from the Unblur corrected micrographs. Initial reference model for 3D reconstruction and refinement was calculated from the atomic structure of archaeal 20S proteasome low-pass filtered to a resolution of 40\AA using `e2pdb2mrc.py` within the EMAN2 package²¹. The resolutions of final 3D reconstructions with D7 symmetry are $2.6(9) \text{\AA}$ calculated from the dataset corrected by Unblur with dose-weighting, $2.5(7)\text{\AA}$ from Unblur followed by particle polishing, $2.5(0)\text{\AA}$ from MotionCor2 with dose-weighting, $2.4(6)\text{\AA}$ from MotionCor2 with particle polishing and $2.4(6)\text{\AA}$ with MotionCor2 with per-frame B-factor weighting respectively (Supplementary Figure 5). Refined maps were low-pass filtered to the nominal resolution reported and sharpened using the automatic B-factor estimation within the RELION post-processing procedure. Various parameters of the different stages of image processing were kept identical between the MotionCor2, and Unblur particle sets such that differences in resulting maps could be attributed primarily to motion correction and dose

weighting.

For the TRPV1 dataset, a total of 236,507 particles were auto-picked from 966 micrographs. All particles were used for reference-based 3D classification implemented in RELION, with C4 symmetry applied. Previously determined TRPV1 density map⁵ was used as the initial model. A total of 71,825 particles from 3D classes with correct structural features were combined for subsequent 3D refinement. During 3D refinement, we found that applying a soft mask to exclude the density for the cytosolic ankyrin domains during the last five refinement iterations slightly improved the final resolution. Resolutions of final 3D reconstruction calculated from the MotionCor2 dose-weighted, and un-weighted, and original MotionCorr particles are 3.15Å, 3.28Å, and 3.43Å, respectively.

For the 20S proteasome dataset collected with 200kV, image motion was corrected using MotionCor2 only. A total of 36,507 particles were auto-picked and screened by 2D classification from 140 micrographs. Final 3D reconstruction with a D7 symmetry has a resolution of 3.48Å. Furthermore, a 3D reconstruction from a subset of 3,297 particles with defocus between 0.5 – 1.0µm has the resolution of 4.08Å, resolving β-strands within β-sheets. Another 3D reconstruction from a different subset of 3,297 particles with defocus between 2.0 – 2.60µm has the resolution of 4.94Å, where β-strands were not resolved.

Model Refinement and Visualization Atomic model of the 20S proteasome⁴ was optimized manually in COOT²² and then automatically with phenix.real_space_refine²³. All map and model visualization was done in UCSF Chimera.

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